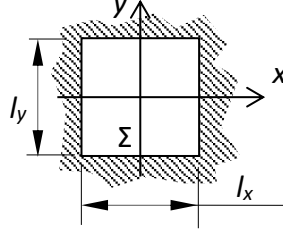


PRACTICE 8

Dr. Gábor Erdei, 19-11-2025

1 Determining the far-field light distribution of a rectangular aperture by Fraunhofer diffraction



$$E(x', y') = \frac{e^{ikz'} \cdot e^{i\frac{k}{2z'}(x'^2 + y'^2)}}{i \cdot \lambda z'} \iint_{\Sigma} E(x, y) \cdot e^{-i\frac{k}{z'}(x'x + y'y)} dx dy \quad (1)$$

$$E(x, y) = \text{rect}\left(\frac{x}{l_x}\right) \cdot \text{rect}\left(\frac{y}{l_y}\right) ; E_0 := 1 ; \text{rect}(a) \triangleq \begin{cases} 1 & \text{if } |a| \leq \frac{1}{2} \\ 0 & \text{if } |a| > \frac{1}{2} \end{cases} \quad (2)$$

$$E(x', y') = \frac{e^{ikz'} \cdot e^{i\frac{k}{2z'}(x'^2 + y'^2)}}{i \cdot \lambda z'} \int_{-l_y/2}^{l_y/2} \int_{-l_x/2}^{l_x/2} e^{-i\frac{k}{z'}(x'x + y'y)} dx dy \quad (3)$$

$$\begin{aligned} \int_{-l_x/2}^{l_x/2} e^{-i\frac{k}{z'}(x'x)} dx &= \left[-\frac{z'}{ikx'} e^{-i\frac{k}{z'}(x'x)} \right]_{-l_x/2}^{l_x/2} = -\frac{z'}{ikx'} \left[e^{-i\frac{k}{z'}(x' \frac{l_x}{2})} - e^{+i\frac{k}{z'}(x' \frac{l_x}{2})} \right] = \\ &= -\frac{z'}{ikx'} \left[-2i \sin \frac{k}{z'} \left(x' \frac{l_x}{2} \right) \right] = \frac{2z'}{kx'} \left[\sin \frac{kx'}{2z'} l_x \right] = \frac{l_x}{l_x} \frac{\lambda z'}{\pi x'} \left[\sin \frac{\pi x' l_x}{\lambda z'} \right] = l_x \text{sinc} \left(\frac{x' l_x}{\lambda z'} \right) \end{aligned} \quad (4)$$

$$\text{sinc}(\xi) \triangleq \frac{\sin \pi \xi}{\pi \xi} ; \quad \xi \triangleq \frac{x' l_x}{\lambda z'} \quad (5)$$

$$E(x', y') = \frac{e^{ikz'} \cdot e^{i\frac{k}{2z'}(x'^2 + y'^2)}}{i \cdot \lambda z'} l_x l_y \text{sinc} \left(\frac{x' l_x}{\lambda z'} \right) \text{sinc} \left(\frac{y' l_y}{\lambda z'} \right) \quad (6)$$

$$I(x', y') = \frac{v \varepsilon}{2} \frac{1}{(\lambda z')^2} l_x^2 l_y^2 \text{sinc}^2 \left(\frac{x' l_x}{\lambda z'} \right) \text{sinc}^2 \left(\frac{y' l_y}{\lambda z'} \right) \quad (7)$$

Where is the first zero point?

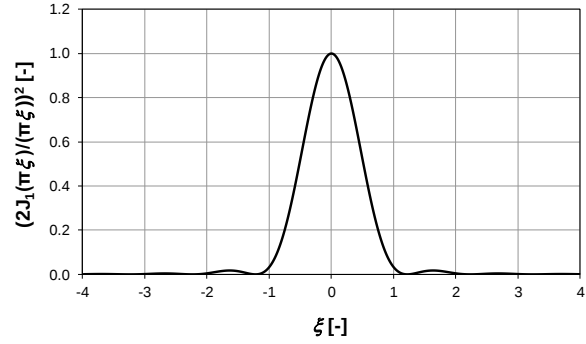
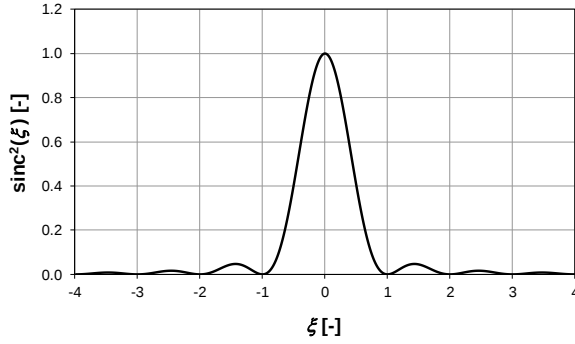
$$\pi \frac{x' l_x}{\lambda z'} = \pi \rightarrow \frac{x' l_x}{\lambda z'} = 1 \quad (8)$$

Nyquist aperture:

$$x_{\text{Nyquist}} \triangleq x' = \frac{\lambda z'}{l_x} \quad (9)$$

In case of a circular aperture (Airy radius):

$$R_{\text{Airy}} \triangleq r' = 1.22 \frac{\lambda z'}{D} ; \quad \xi \triangleq \frac{r' D}{\lambda z'} \quad (10)$$



E.g. What area of the objective lens ($f = 30$ mm) of a video projector is illuminated by one pixel of the spatial light modulator (SLM) of the projector ($l_x = 7.6$ μm), if $\lambda = 550$ nm? The lens is considered as ideal, placed at about a focal distance from the SLM.

$$x_{\text{Nyquist}} = \frac{\lambda z'}{l_x} = \frac{0.550 \cdot 30}{7.6} = 2.17 \text{ mm} \quad (11)$$

2 Resolving power of a telescope, see lecture notes

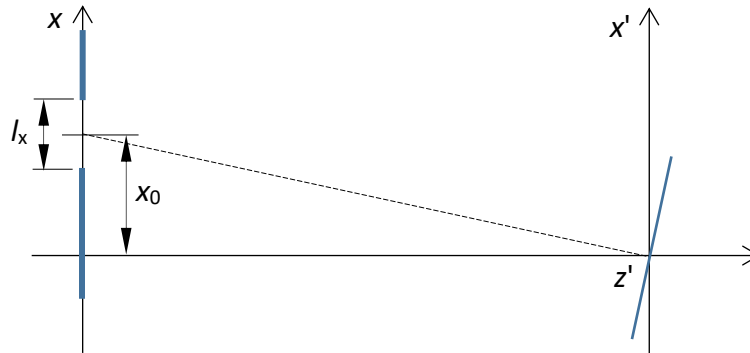
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E.g. Very Large Telescope – VLT (ESO, Paranal Observatory), $D = 8.2$ m, $\lambda = 550$ nm

$$\Delta\theta' \approx \text{tg } \Delta\theta' = \frac{R_{\text{Airy}}}{z'} = 1.22 \frac{\lambda}{D} = 8.2 \cdot 10^{-8} \text{ rad} \quad (12)$$

The angular field-of-view of the largest moon of Saturn, the Titan is: $3.9 \cdot 10^{-6}$ rad!

3 Far-field modelling of a rectangular aperture (slit) shifted from the axis



Let the value of the position shift be x_0 : $E(x, y) \rightarrow E(x - x_0, y)$. According to (4):

$$\begin{aligned} \int_{x_0 - l_x/2}^{x_0 + l_x/2} e^{-i \frac{k}{z'} (x' x)} dx &= \left[-\frac{z'}{ikx'} e^{-i \frac{k}{z'} (x' x)} \right]_{x_0 - l_x/2}^{x_0 + l_x/2} = \\ &= -\frac{z'}{ikx'} e^{-i \frac{k}{z'} (x' x_0)} \left[e^{-i \frac{k}{z'} (x' \frac{l_x}{2})} - e^{+i \frac{k}{z'} (x' \frac{l_x}{2})} \right] = \\ &= e^{-i \frac{kx_0}{z'} x'} \cdot l_x \text{sinc} \left(\frac{x' l_x}{\lambda z'} \right) \end{aligned} \quad (13)$$

In the obtained result a plane wave appears, which exactly corresponds to the shift theorem of Fourier transformation, if we take the value of spatial frequency to be $f_x = x'/(\lambda z')$. The shift of the slit does not cause any changes to the intensity distribution.

4 Far-field modelling of two rect. apertures shifted symmetrically from the axis

Let the shift be $+x_0$ and $-x_0$: $E(x, y) \rightarrow E(x - x_0, y) + E(x + x_0, y)$. The calculation makes sense only when the slits do not overlap, i.e. $x_0 \geq l_x/2$. If $l_y \gg l_x$, then the structure can be considered as one-dimensional. Based on (13):

$$\int_{x_0 - l_x/2}^{x_0 + l_x/2} e^{-i\frac{k}{z'}(x'x)} dx + \int_{-x_0 - l_x/2}^{-x_0 + l_x/2} e^{-i\frac{k}{z'}(x'x)} dx =$$

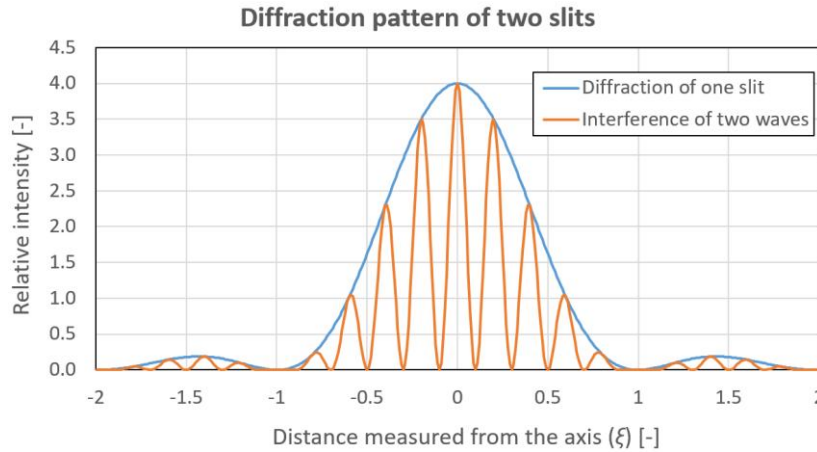
$$\left(e^{-i\frac{kx_0}{z'}x'} + e^{+i\frac{kx_0}{z'}x'} \right) \cdot l_x \operatorname{sinc}\left(\frac{x'l_x}{\lambda z'}\right) = 2 \cos\left(\frac{kx_0}{z'}x'\right) \cdot l_x \operatorname{sinc}\left(\frac{x'l_x}{\lambda z'}\right) \quad (14)$$

So, adding a second slit does change the intensity distribution too:

$$I(x', y') = \frac{v\epsilon}{2} \frac{1}{(\lambda z')^2} l_x^2 l_y^2 \operatorname{sinc}^2\left(\frac{x'l_x}{\lambda z'}\right) \operatorname{sinc}^2\left(\frac{y'l_y}{\lambda z'}\right) 4 \cos^2\left(\pi \frac{x'2x_0}{\lambda z'}\right) \quad (15)$$

Where $2x_0 = d$ is the distance between the slits. The above relationship is the scalar diffraction model of Young's two-slit experiment. According to $2 \cos^2(a) \equiv 1 + \cos(2a)$ we even get back the interferogram of two oblique plane waves (see lecture 7). The relationship is in perfect correspondence with the expression obtained at two-beam interference:

$$4 \cos^2\left(\pi \frac{x'2x_0}{\lambda z'}\right) = 2 \left(1 + \cos\left(2\pi \frac{x_0}{\lambda} \frac{x'}{z'}\right) \right) = 2(1 + \cos(\Delta k' x')) \quad (16)$$



$d/l_x = 5$.

E.g. Let us place a single hair (of 50 μm thickness) in the middle of a narrow, long slit of 150 μm thickness, parallel with its longer side. The optical model of the arrangement is two slits of $l_x = 50 \mu\text{m}$ width, with their centers placed $d = 100 \mu\text{m}$ apart. Since $l_y \gg l_x$ we make the calculation in one dimension, $\lambda = 633 \text{ nm}$. What is the distance of the interference fringes on a screen being at a distance of $z' = 200 \text{ mm}$?

$$2 \frac{2\pi x_0}{\lambda} \frac{x'}{z'} = 2\pi \rightarrow x' = \frac{\lambda z'}{2x_0} = \frac{\lambda z'}{d} \quad (17)$$

$$x' = \frac{\lambda z'}{d} = \frac{0.633 \mu\text{m} \cdot 200 \text{ mm}}{100 \mu\text{m}} = 1.27 \text{ mm} \quad (18)$$

How many zeros does the interferogram have between the first zero points of the diffraction pattern of one slit (i.e. inside the two-times Nyquist aperture)?

$$2x_{\text{Nyquist}} = 2 \frac{\lambda z'}{l_x} \quad (19)$$

$$\frac{2x_{\text{Nyquist}}}{x'} = \frac{2d}{l_x} = 4 \text{ pcs.} \quad (20)$$